

# A Heterogeneous Agent Model of Asset Price with Three Time Delays

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This paper considers a continuous-time heterogeneous agent model of a financial market with one risky asset, two types of agents (i.e., the fundamentalists and the chartists), and three time delays. The chartist's demand is determined through a nonlinear function of the difference between the current price and a weighted moving average of the delayed prices whereas the fundamentalist's demand is governed by the difference between the current price and the fundamental value. The asset price dynamics is described by a nonlinear delay differential equation. Two main results are analytically and numerically shown:

 (i) a single delay destabilizes the market price and generates cyclic oscillations around the equilibrium;

(i) under multiple delays, stability loss and gain repeatedly occur as the length of the

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# **1. INTRODUCTION**

delay increases.

Persistent volatility is a prominent characteristic feature of financial markets. It is, however, well-known that the efficient market hypothesis cannot account for the discrepancy between the observed market price and the fundamental value of the asset. Recently, heterogeneous agent models (HAMs) have been developed to explain a wide variety of financial market behavior such as temporary bubbles, sudden market crashes and price resistance in discrete-time as well as in continuous-time framework. In his survey of recent developments of HAMs, Hommes [1] discusses that nonlinear discrete-time HAMs can generate various dynamics ranging from cyclic fluctuations to chaotic behavior. Among others there is Chiarella et al. [2] who propose a discrete-time HAM with a moving average (MA) rule having various memory lengths. Their main finding is that the length of the MA rule can be a source of complicated dynamics in destabilized financial markets in deterministic and stochastic processes of difference equations. On the other hand, continuoustime HAMs are also examined in various ways in which dynamics are described by ordinary or delay differential equations. They have a long history since Zeeman [3] and Beja and Goldman [4]. Moreover, Chiarella [5], which is a development of Beja and Goldman [4], shows that the market price tends to a stable limit cycle under a nonlinear demand function of the risky asset when the equilibrium is unstable. More recently, He and Zheng [6] reconstruct the discretetime model of Chiarella et al. [2] in a continuous-time framework in which the expected price is formed with a moving average of the past (delay) prices. Their main result concerns a double edge effect on the stability caused by the length of the memory or delay: an increase in delay can destabilize the market price and also stabilize it. In their model as well as in subsequent studies on

continuous-time HAMs such as He and Li [7] and Xu et al. [8], this interesting phenomenon is obtained under the assumption that infinitely many past price data are available at no charge of cost. Needless to remember the well-known line in economics, "there is no such thing as a free lunch," it is to go too far to get necessary information for nothing. At least two ways are possible to render this extreme assumption to more realistic one, the first is to introduce information cost associated with collecting the price data and the second is to limit availability of the past prices. In the present paper, we take the second way and then reconsider the price dynamics when the chartists forecast an expected price, using only a limited information on past prices.

Following the framework of Dibeh [9], this paper constructs a dynamic HAM of a speculative asset with three delays (i.e., three past prices) and investigates the effect of time delays on the asset price dynamics. Dibeh [9] conducts mainly with numerical simulations. Analytical developments of a one-delay version and a two-delay version of Dibeh's model are already presented by Qu and Wei [10] and Matsumoto and Szidarovszky [11], respectively. The model we analyze is a continuation of these preceding studies. One of our main findings is that stability switches from stability to instability as well as from instability to stability can occur in three delay framework. That is, the double edge effect can also be observed in continuous-time HAMs with a finite number of delay prices.

The structure of the paper is as follows. Section Model constructs a HAM model with fundamentalists and chartists. The stability switching curve is analytically derived, which divides the delay space into stability and instability regions. Section Dynamics has three subsections, in each of which we conduct a stability and bifurcation analysis under different circumstances. Section Concluding Remarks concludes the paper with further research directions.

## 2. MODEL

We consider an asset pricing model with one risky asset and two traders, fundamentalists and chartists. Let p(t) be the price of the risky asset at time *t*. Since the fundamentalists believe that p(t) eventually converges to *f*, the fundamental (i.e., equilibrium) price of the asset, they will sell the asset if p(t) > f and buy it if p(t) < f. The simplest form of the demand function of the fundamentalist is

$$D^{f}(p(t)) = m\left[f - p(t)\right]$$

where  $m \in (0, 1)$  is the fraction of the fundamentalists in the market. On the other hand, the chartists base their decisions of market participation on the price trend of the asset. Their demand function is

$$D^{c}(s(t)) = (1 - m)g(s(t))$$

where s(t) is a weighted average of the past price trends formulated at time t,

$$s(t) = \sum_{i=1}^{n} \alpha_i \left[ p(t - \tau_{i-1}) - p(t - \tau_i) \right]$$

where  $\tau_i \ge 0$  is a time delay with  $\tau_0 = 0$  and  $\alpha_i \ge 0$  denotes a weight of a price change satisfying  $\sum_{i=1}^{n} \alpha_i = 1$ . The demand function g(s) is assumed to be hyperbolic tangent as in Chiarella [5] and Dibeh [9],

$$g(s) = \tanh(s).$$

The market demand is the sum of demands of the fundamentalists and chartists,

$$D(p(t), s(t)) = D^{c}(s(t)) + D^{f}(p(t)).$$

The price dynamics follows Dibeh's formulation in which the growth rate of price is determined by the market demand,

$$\frac{\dot{p}(t)}{p(t)} = D(p(t), s(t))$$

that is written as a nonlinear delay differential equation,

$$\dot{p}(t) = (1 - m)p(t) \tanh[s(t)] + mp(t)[f - p(t)].$$
 (1)

Notice that s(t) can be rewritten as

$$s(t) = \alpha_1 \left[ p(t) - \bar{p}_n(t) \right]$$

where

$$\bar{p}_n(t) = \sum_{j=1}^n \eta_j p(t - \tau_j)$$

with

$$\eta_j = \frac{\alpha_j - \alpha_{j+1}}{\alpha_1}$$
 for  $1 \le j \le n-1$  and  $\eta_n = \frac{\alpha_n}{\alpha_1} \left( = 1 - \sum_{j=1}^{n-1} \eta_j \right)$ .

Since  $\bar{p}_n(t)$  is considered to be the weighted (moving) average of the past *n* prices, the chartists believe that the price will rise (fall) when the current price is above (below) the average of the *n* delay prices. If infinitely many price data is available with finite or infinite memory length, the moving average can be presented by continuously distributed time delay such as

$$p_{\xi}(t) = \int_{t-\xi}^{t} \eta(s) p(s) ds$$

where  $0 < \xi \leq \infty$  is the memory length and  $\eta(s)$  is a weighting function. This form is used in He and Zheng [6].

Concerning the specification of s(t), Qu and Wei [10] examine the case of n = 1 in which the trend includes one delay price with  $\alpha_1 = 1$ 

$$s(t) = p(t) - p(t - \tau_1).$$

It is shown first that there exists a threshold value of the delay at which stability is lost and second that the stationary point can be stable for smaller values and bifurcates to a limit cycle for larger values. Matsumoto and Szidarovszky [11] consider the case of n = 2 in which two delay prices are used to formulate the trend,

$$s(t) = \alpha_1 \left( p(t) - p(t-\tau_1) \right) + \alpha_2 \left( p(t-\tau_1) - p(t-\tau_2) \right).$$

or

$$S(t) = \alpha_1 \left\{ p(t) - \left[ (1 - \eta)p(t - \tau_1) + \eta p(t - \tau_2) \right] \right\}$$

with

5

$$\eta = \frac{\alpha_2}{\alpha_1}.$$

With two delay prices, the coefficients are positive and add up to unity in interpolation (i.e.,  $\alpha_1 > \alpha_2$ ) whereas one coefficient is negative, the other greater than unity and the sum is also unity in extrapolation (i.e.,  $\alpha_1 < \alpha_2$ ). It is demonstrated that the two delay model has a double edge effect with which stability losses and gains can repeatedly occur.

In this study, we move one more step forward and draw attention to the case of n = 3 for which the average trend is

$$s(t) = \alpha_1 [p(t) - p(t - \tau_1)] + \alpha_2 [p(t - \tau_1) - p(t - \tau_2)] + \alpha_3 [p(t - \tau_2) - p(t - \tau_3)]$$

or

$$s(t) = \alpha_1 \left\{ p(t) - \left[ \eta_1 p(t - \tau_1) + \eta_2 p(t - \tau_2) + \eta_3 p(t - \tau_3) \right] \right\}$$

with

$$\eta_1 = \frac{\alpha_1 - \alpha_2}{\alpha_1}, \ \eta_2 = \frac{\alpha_2 - \alpha_3}{\alpha_1} \ \text{and} \ \eta_3 = \frac{\alpha_3}{\alpha_1}.$$

For the sake of convenience, we adopt the following form,

$$s(t) = \alpha_1 p(t) + (\alpha_2 - \alpha_1) p(t - \tau_1) + (1 - \alpha_1 - 2\alpha_2) p(t - \tau_2) -(1 - \alpha_1 - \alpha_2) p(t - \tau_3).$$

The unique stationary point of dynamic Equation (1) is identical with the equilibrium price,

$$p^e = f = p(t) = p(t - \tau_j)$$
 for all  $t \ge 0$  and  $j = 1, 2, 3$ .

It is clear that at the stationary point there is no price trend,  $s^e = 0$ . Linearizing the nonlinear dynamic equation in a neighborhood of the stationary point yields

$$\dot{p}_{\delta}(t) = \alpha p_{\delta}(t) + \beta_1 p_{\delta}(t - \tau_1) + \beta_2 p_{\delta}(t - \tau_2) - \beta_3 p_{\delta}(t - \tau_3)$$
(2)

where  $p_{\delta}(t) = p(t) - p^e$  and the coefficients are defined as

$$\begin{aligned} \alpha &= f \left[ (1 - m)\alpha_1 - m \right], \\ \beta_1 &= f (1 - m)(\alpha_2 - \alpha_1), \\ \beta_2 &= f (1 - m)(1 - \alpha_1 - 2\alpha_2) \\ \beta_3 &= f (1 - m)(1 - \alpha_1 - \alpha_2). \end{aligned}$$





With condition 
$$\sum_{i=1}^{3} \alpha_i = 1$$
, it can be verified that

$$\begin{array}{l} \alpha_3 = 0 & \Longleftrightarrow & \alpha_2 = 1 - \alpha_1, \\ \alpha_3 = \alpha_1 & \Longleftrightarrow & \alpha_2 = 1 - 2\alpha_1, \\ \alpha_3 = \alpha_2 & \Longleftrightarrow & \alpha_2 = \frac{1}{2}(1 - \alpha_1). \end{array}$$

The locus of  $\alpha_3 = 0$  divides the first quadrant of the  $(\alpha_1, \alpha_2)$  plane into two parts, upper and lower right triangles, and condition  $\alpha_3 \ge 0$  eliminates the upper one. The lower triangle is further divided into six subparts by the three loci of  $\alpha_2 = \alpha_1$ ,  $\alpha_3 = \alpha_2$ , and  $\alpha_3 = \alpha_1$ . Those divisions are depicted in **Figure 1** in which the relations among the magnitudes of  $\alpha_1$ ,  $\alpha_2$ , and  $\alpha_3$  are determined in the following way,

- $\alpha_1 > \alpha_2 > \alpha_3$  in region I,  $\alpha_3 > \alpha_2 > \alpha_1$  in region IV,
- $\alpha_1 > \alpha_3 > \alpha_2$  in region II,  $\alpha_2 > \alpha_3 > \alpha_1$  in region V,
- $\alpha_3 > \alpha_1 > \alpha_2$  in region III,  $\alpha_2 > \alpha_1 > \alpha_3$  in region VI.

Region I is colored in gray and we will limit our analysis to it below. The line of  $\alpha_2 = (2 - 3m)/2(1 - m) - \alpha_1$  is shown to be dotted and downward sloping<sup>1</sup>. It is also checked that  $\beta_1 = 0$ ,  $\beta_2 = 0$ , and  $\beta_3 = 0$  hold, respectively, on the loci of  $\alpha_2 = \alpha_1$ ,  $\alpha_3 = \alpha_2$ , and  $\alpha_3 = 0$ . On these boundaries the threedelay Equation (2) is reduced to one of three two-delay equations, depending on which  $\beta_i$  value becomes zero. Matsumoto and Szidarovszky [11] consider a two delay equation that corresponds to the one with  $\beta_3 = 0$ . Their analytical method can be applied to the other two equations as well. In this study, we will confine

 $<sup>^{1}</sup>m = 0.4$  is assumed and we will refer to this line later.

our attention to the regions in which  $\beta_i > 0$ . Although, the parametric region I is a smaller part of the whole feasible region, the three delay model restricted to region *I* has more freedom concening the parametric choice than the two delay model that is defined only on the line  $\alpha_2 = 1 - \alpha_1$ .

# 2.1. Construction of Stability Switching Curves

Substituting an exponential solution  $p_{\delta}(t) = e^{\lambda t} u$  into the linearized Equation (2) yields the corresponding characteristic equation

$$\lambda - \alpha - \beta_1 e^{-\lambda \tau_1} - \beta_2 e^{-\lambda \tau_2} + \beta_3 e^{-\lambda \tau_3} = 0$$
 (3)

Dividing the left hand side of Equation (3) by  $\lambda - \alpha$  and denote the result by  $a(\lambda)$ ,

$$a(\lambda) = 1 + a_1(\lambda)e^{-\lambda\tau_1} + a_2(\lambda)e^{-\lambda\tau_2} + a_3(\lambda)e^{-\lambda\tau_3} = 0$$
 (4)

with

$$a_1(\lambda) = \frac{-\beta_1}{\lambda - \alpha}, a_2(\lambda) = \frac{-\beta_2}{\lambda - \alpha}, a_3(\lambda) = \frac{\beta_3}{\lambda - \alpha}.$$

Since  $\lambda = 0$  is not a solution of Equation (3), a pair of pure imaginary roots must exist if stability switch occurs. We thus assume  $\lambda = i\omega$ ,  $\omega > 0$  is a solution of Equation (3). For  $\lambda = i\omega$ ,

$$a_1(i\omega) = \frac{\alpha\beta_1 + i\beta_1\omega}{\alpha^2 + \omega^2} \text{ and } |a_1(i\omega)| = \frac{|\beta_1|}{\sqrt{\alpha^2 + \omega^2}},$$
$$a_2(i\omega) = \frac{\alpha\beta_2 + i\beta_2\omega}{\alpha^2 + \omega^2} \text{ and } |a_2(i\omega)| = \frac{|\beta_2|}{\sqrt{\alpha^2 + \omega^2}}$$

and

$$a_3(i\omega) = -\frac{\alpha\beta_3 + i\beta_3\omega}{\alpha^2 + \omega^2}$$
 and  $|a_1(i\omega)| = \frac{\beta_3}{\sqrt{\alpha^2 + \omega^2}}$ 

Each term of  $a(i\omega)$  may be viewed as a vector in the complex plane. Hence solving  $a(i\omega) = 0$  analytically is equivalent to constructing a quadrangle geometrically as shown in **Figure 2** in which these four vectors are arranged from head to tail.

The determination of the stability switching surface in the  $(\tau_1, \tau_2, \tau_3)$  plane is challenging<sup>2</sup>. In order to make the problem manageable, we reduce the three-delay equation to a two-delay equation by fixing the value of  $\tau_3$  at a certain positive level and then construct the stability switching curve in the  $(\tau_1, \tau_2)$  plane in which we can apply the method used earlier in Matsumoto and Szidarovszky [11], and Gu et al. [14]. Notice that vector 1 is connected to vector  $a_3(i\omega)e^{-i\omega\tau_3}$  in **Figure 2**. If the sum of these vectors is denoted by  $a_d(i\omega, \tau_3)$ ,

$$a_d(i\omega,\tau_3) = 1 + a_3(i\omega)e^{-i\omega\tau_3},$$

then  $a(\lambda)$  with  $\lambda = i\omega$  can be rewritten as

$$a(i\omega) = a_d(i\omega, \tau_3) + a_1(i\omega)e^{-i\omega\tau_1} + a_2(i\omega)e^{-i\omega\tau_2}$$

<sup>2</sup>It is possible to have a switching surface in the 3D space. See Almodaresi and Bozorg [12] and Gu and Naghnaeian [13].



As already shown in **Figure 2**,  $a(i\omega) = 0$  means that these three vectors in  $a(i\omega)$  must form a triangle having the dotted base with interior angles,  $\theta_1$  and  $\theta_2$ . Sufficient and necessary conditions for forming a triangle are given by

$$(i) f_d(\omega) = |a_1(i\omega)| + |a_2(i\omega)| - |a_d(i\omega, \tau_3)| \ge 0,$$
  

$$(ii) f_1(\omega) = |a_2(i\omega)| + |a_d(i\omega, \tau_3)| - |a_1(i\omega)| \ge 0,$$
 (5)  

$$(iii) f_2(\omega) = |a_d(i\omega, \tau_3)| + |a_1(i\omega)| - |a_2(i\omega)| \ge 0.$$

Each inequality condition implies that the length of any segment of the triangle is not greater than the sum of the lengths of the remaining two segments. Define the crossing frequency set  $\Omega$  of all  $\omega > 0$  such that  $a(i\omega) = 0$  holds for at least one delay combination  $(\tau_1, \tau_2, \tau_3) \ge 0$ . For  $\omega = 0$  in  $a(i\omega)$ leads to

$$a(i\omega)\big|_{\omega=0}=\frac{1}{\alpha}\left(\alpha+\beta_1+\beta_2-\beta_3\right)<0.$$

This inequality implies  $0 \notin \Omega$ .

By the law of cosine, we can determine the values of  $\theta_1$  and  $\theta_2$  of the triangle in **Figure 2**,

$$\theta_{1}(\omega) = \cos^{-1} \left[ \frac{|a_{d}(i\omega, \tau_{3})|^{2} + |a_{1}(i\omega)|^{2} - |a_{2}(i\omega)|^{2}}{2 |a_{d}(i\omega, \tau_{3})| |a_{1}(i\omega)|} \right]$$

and

$$\theta_2(\omega) = \cos^{-1}\left[\frac{\left|a_d(i\omega, \tau_3)\right|^2 + \left|a_2(i\omega)\right|^2 - \left|a_1(i\omega)\right|^2}{2\left|a_d(i\omega, \tau_3)\right| \left|a_2(i\omega)\right|}\right]$$

Vertices *A* and *B* maybe located above or below the horizontal axis and the slope of segment *AB* can be positive or negative. So

we have eight possibilities to construct a quadrangle. However, a simple geometric consideration shows that there are only two different possibilities:

$$\arg\left(a_1(i\omega)e^{-i\omega\tau_1}\right) + (2k-1)\pi - \arg(a_d(i\omega,\tau_1)) \pm \theta_1(\omega) = 0$$

and

$$\arg\left(a_2(i\omega)e^{-i\omega\tau_2}\right) + (2n-1)\pi - \arg(a_d(i\omega,\tau_1)) \mp \theta_2(\omega) = 0.$$

Solving these equations for  $\tau_1$  and  $\tau_2$  yields the threshold values of the delays,

$$\tau_1^{\pm}(\omega, k) = \frac{1}{\omega} \left[ \arg(a_1(i\omega)) - \arg(a_d(i\omega, \tau_3)) + (2k - 1)\pi \pm \theta_1(\omega) \right]$$
(6)

and

$$\tau_2^{\mp}(\omega, n) = \frac{1}{\omega} \left[ \arg(a_2(i\omega)) - \arg(a_d(i\omega, \tau_3)) + (2n-1)\pi \\ \mp \theta_2(\omega) \right].$$
(7)

The stability switching curves consist of two sets of parametric segments,

$$L_1(k, n) = \{\tau_1^+(\omega, k), \tau_2^-(\omega, n)\} \text{ for } k, n = 0, 1, 2, \dots$$
(8)

and

$$L_2(k, n) = \{\tau_1^-(\omega, k), \tau_2^+(\omega, n)\} \text{ for } k, n = 0, 1, 2, \dots$$
(9)

where the parameter  $\omega$  runs through the crossing frequency set  $\Omega$ .

Substituting  $\tau_1 = \tau_2 = \tau_3 = 0$  into Equation (3) presents

$$\lambda = \alpha + \beta_1 + \beta_2 - \beta_3 = -fm < 0.$$

This inequality implies that the steady state with no delays is stable. Next, assuming that  $\tau_1 = \tau_2 = 0$ , we examine the existence of a threshold value of  $\tau_3$  at which the system loses stability. Without loss of generality, assuming  $\lambda = i\omega$ ,  $\omega > 0$  and then substituting it into Equation (3) with  $\tau_1 = \tau_2 = 0$  reduce the characteristic equation to

$$i\omega - (\alpha + \beta_1 + \beta_2) + \beta_3 e^{-i\omega\tau_3} = 0.$$

We denote  $\alpha + \beta_1 + \beta_2$  by  $\Delta$  where

$$\Delta = f(1-m)\left(1-\alpha_1-\alpha_2-\frac{m}{1-m}\right)$$

and break down the characteristic equation into the real and imaginary parts,

$$-\Delta + \beta_3 \cos \omega \tau_3 = 0$$
(10)  
$$\omega - \beta_3 \sin \omega \tau_3 = 0.$$

Moving  $\Delta$  and  $\omega$  to the right hand side and adding the squared equations give

$$\omega^2 = \beta_3^2 - \Delta^2$$

 $\beta_3 - \Delta = fm > 0$ 

 $\beta_3 + \Delta = 2f(1-m)\left(\varphi(\alpha_1) - \alpha_2\right)$ 

where

and

with

$$\varphi(\alpha_1) = \frac{2 - 3m}{2(1 - m)} - \alpha_1.$$

Notice that the dotted curve in **Figure 1** is described by  $\alpha_2 = \varphi(\alpha_1)$ . It is clear first that

$$\beta_3 + \Delta \gtrless 0$$
 according to  $\varphi(\alpha_1) \gtrless \alpha_2$ 

and second that

$$\beta_3 + \Delta < 0$$
 if  $m \ge 2/3$ .

This leads to the following result.

Proposition 1. Given  $\tau_1 = \tau_2 = 0$ , the solution of delay system (1) is stable for any  $\tau_3 \ge 0$  if the fundamentalists dominate over the chartists in the sense that  $m \ge 2/3$ .

If  $\beta_3 + \Delta > 0$ , then there is a  $\omega$  such as  $\bar{\omega} = \sqrt{\beta_3^2 - \Delta^2} > 0$ . In this case, loss of stability for  $\lambda = i\bar{\omega}$  can be shown in the following way. The characteristic equation with  $\tau_1 = \tau_2 = 0$  can be written as

$$\beta_3 e^{-\lambda \tau_3} = \beta_1 + \beta_2 + \alpha - \lambda. \tag{11}$$

Taking  $\lambda$  as a function of  $\tau_3$  and differentiating it with respect to  $\tau_3$  yield

$$\frac{d\lambda}{d\tau_3} = \frac{\lambda\beta_3 e^{-\lambda\tau_3}}{1 - \tau_3\beta_3 e^{-\lambda\tau_3}}$$
$$= \frac{\lambda(\beta_1 + \beta_2 + \alpha - \lambda)}{1 - \tau_3(\beta_1 + \beta_2 + \alpha - \lambda)}$$

where the right hand side of Equation (11) are used. Substituting  $\lambda = i\omega$  and taking the real part of the resulting expression give

$$\frac{d[\operatorname{Re}\lambda]}{d\tau_{3}}\Big|_{\lambda=i\omega} = \operatorname{Re}\left[\frac{\lambda(\beta_{1}+\beta_{2}+\alpha-\lambda)}{1-\tau_{3}(\beta_{1}+\beta_{2}+\alpha-\lambda)}\right] = \operatorname{Re}\left[\frac{\omega^{2}+i\omega(\beta_{1}+\beta_{2})}{1-\tau_{3}(\beta_{1}+\beta_{2}+\alpha)-i\omega}\frac{1+\tau_{3}(\beta_{1}+\beta_{2}+\alpha)-i\omega}{1+\tau_{3}(\beta_{1}+\beta_{2}+\alpha)-i\omega}\right] = \frac{\omega^{2}}{\left[1-\tau_{3}(\beta_{1}+\beta_{2}+\alpha)\right]^{2}+\omega^{2}} > 0$$

This inequality implies that all roots cross the imaginary axis at  $i\omega$  from left to right as  $\tau_3$  increases so stability is lost. Correspondingly, solving the first equation of Equation (10) for  $\tau_3$  presents the threshold value

$$\bar{\tau}_3(\alpha_1, \alpha_2, m) = \frac{1}{\bar{\omega}} \cos^{-1}\left(\frac{\Delta}{\beta_3}\right) > 0 \tag{12}$$

and the system is stable for  $\tau_3 < \bar{\tau}_3$  and unstable otherwise<sup>3</sup>. We summarize the results obtained so far.

Proposition 2. Given  $\tau_1 = \tau_2 = 0$ , the solution of delay system (1) is stable for  $\tau_3 < \overline{\tau}_3(\alpha_1, \alpha_2, m)$ , loses stability for  $\tau_3 = \overline{\tau}_3(\alpha_1, \alpha_2, m)$  and becomes unstable for all  $\tau_3 > \overline{\tau}_3(\alpha_1, \alpha_2, m)$  if the following condition holds,

$$\alpha_2 < \varphi(\alpha_1)$$

whereas it is always stable for any  $\tau_3 > 0$  otherwise.

It is to be noticed that with  $\tau_1 = \tau_2 = 0$ , the three delay model is reduced to a one delay model and thus this stability result is essentially the same as the one of Qu and Wei [10]. In the next section we will proceed to the case of  $\tau_i > 0$  for i = 1, 2, 3.

# 3. DYNAMICS

Central to this study is the problem of dynamics associated with three fixed delays. This problem is complex and it seems unlikely that an analytical approach to nonlinear delay differential Equation (1) would generate fruitful results on dynamics as it might be difficult to solve the equation. In order to understand global dynamics with delays better, it is useful to take a numerical approach. To this end, we make several simplifying assumptions and then numerically examine dynamics of the asset price.

**Assumption 1**. f = 5, m = 0.4 and  $\tau_3 = 5$ .

**Assumption 2.**  $\alpha_i$  for i = 1, 2, 3 satisfy  $\alpha_1 > \alpha_2 > \alpha_3 > 0$ .

**Assumption 3**.  $\tau_i$  for i = 1, 2, 3 satisfy  $\tau_1 < \tau_2 < \tau_3$ .

Since the parameter values in Assumption 1 are determined only for computational convenience, the main results to be obtained below hold for any other values, provided that m < 2/3. Assumption 2 implies that the more recent trend has the more weight. This is not necessary but simplifies the analysis. Under m = 0.4 of Assumption 1, region I is located above the line  $\alpha_2 = \varphi(\alpha_1)$  which is the dotted downward-sloping line passing through point (1/3, 1/3) in **Figure 1**. Due to Proposition 2, the stationary point is always stable in region I where  $\alpha_2 > \varphi(\alpha_1)$  if  $\tau_1 = \tau_2 = 0$ . Thus it is easy to see how increasing values of these delays affect stability of the steady state. Assumption 3 imposes a natural ordering among the three delays, that is, an older price has a longer delay.

The rest of this section is divided into three subsections. We will look at the  $\tau_1$ -effect on dynamics caused by a change in  $\tau_1$  in Section Delay Effect I:  $\tau_1$ -effect, explore the  $\tau_2$ -effect in Section Delay Effect II:  $\tau_2$ -effect, and focus on considerations of the  $\tau_3$ -effect in Section Delay Effect III:  $\tau_3$ -effect.

## **3.1. Delay Effect I:** $\tau_1$ -effect

We start with the  $\tau_1$ -effect and present some numerical examples to see how different values of  $\tau_1$  affect dynamics of Equation (1). We first check the triangle conditions. Subtracting the third equation from the second equation Equation in (5) yields

$$f_1(\omega) - f_2(\omega) = \frac{6f(1-m)}{\sqrt{\alpha^2 + \omega^2}} \left(\alpha_2 - \frac{1}{3}\right)$$
 (13)

implying that the relative location of the curves of  $f_1(\omega)$  and  $f_2(\omega)$  depends on whether the value of  $\alpha_2$  is larger or less than 1/3. To simplify the analysis, we assume that the value of  $\alpha_2$  is kept constant at 1/3;

Assumption 4.  $\alpha_2 = 1/3$ 

Assumption 2 implies  $\eta_i > 0$  for all *i*, that is, the weights of the delay prices to calculate the moving average are positive. Their relative magnitudes depend on the selected values of  $\alpha_i$ and thus there are many possibilities. Among them, we will limit analysis to only two cases,  $\eta_1 = \eta_2 > \eta_3$  if  $\alpha_1 < 1/2$  and  $\eta_3 > \eta_1 = \eta_2$  if  $\alpha_1 < 1/2$  since Assumption 4 makes  $\eta_1 = \eta_2$ always<sup>4</sup>. With Assumption 4,  $f_1(\omega) = f_2(\omega)$  holds regardless of the value of  $\alpha_1$ . Further,  $f_i(\omega) > 0$ , i = 1, 2, for  $\omega > 0$ which will be numerically confirmed below. Consequently, it is enough to verify only the location and the shapes of the  $f_d(\omega)$ curve. We consider dynamics over an interval  $[\alpha_1^0, 2/3]$  of  $\alpha_1$ . The left hand side extreme value,  $\alpha_1^0 \simeq 0.356$  is numerically determined so as to make the maximum of  $f_d(\omega)$  equal to zero. Thus for  $\alpha_1 < \alpha_1^0$ ,  $f_d(\omega) < 0$  for all  $\omega > 0$ , implying that one of the triangle conditions is always violated. The other extreme value,  $\alpha_1 = 2/3$ , together with  $\alpha_2 = 1/3$  makes  $\alpha_3 = 0$  via  $\alpha_2 = 1 - \alpha_1$ . For  $\alpha_1 > 2/3$ , the nonnegative condition  $\alpha_3 \ge 0$ is violated. In the following numerical analysis,  $\alpha_1$  is increased from 6/15 to 9/15 with an increment of 1/15. The resultant four  $f_d(\omega)$  curves are colored in black in **Figure 3** where the red curve with  $\alpha_1 = 1/2$  and the blue curves with those extreme values are also depicted. The corresponding  $f_1(\omega) = f_2(\omega)$  curves are illustrated in the dotted curves with the same color and are seen to be positive for all  $\omega \geq 0$ . In the gray region below the lower blue curve, no stability switch occurs and thus the steady state is stable. We do not consider the  $\tau_1$ -effect there. On the other hand, the gray region above the upper blue curve is not feasible and thus eliminated from consideration. Figure 3 also illustrates various shapes of the  $f_d(\omega)$  curves and shows that increasing  $\alpha_1$ shifts the curve upward and removes its unevenness. Indeed, the lower blue curve with  $\alpha_1 = \alpha_1^0$  located at the bottom of the yellow region is high-wave shaped while the higher blue curve with  $\alpha_1 = 2/3$  at the top is monotonically downward-sloping. In addition, as the red curve passes through the origin, we then see that the curves below the red curve have two intersections with the horizontal axis and the curves above have only one intersection.

To be more specific, we pick up  $\alpha_1 = 6/15$  and construct a stability switching curve. The corresponding  $f_d(\omega)$  curve is

 $<sup>^3</sup>$  Solving the second equation of Equation (10) for  $\tau_3$  gives the same result in a different form.

<sup>&</sup>lt;sup>4</sup>The remaining cases will be examined in a subsequent paper.

the black one just above the lower blue curve in **Figure 3** and intersects the horizontal axis twice at

$$ar{\omega}_1\simeq 0.438$$
 and  $ar{\omega}_2\simeq 0.582$ 

The triangle conditions in Equation (5) hold for  $\omega \in \Omega = [\bar{\omega}_1, \bar{\omega}_2]$ . We now ready to derive the stability switching curves,  $L_1(k, n)$  and  $L_2(k, n)$  defined in (8) and (9). Taking k = n = 1, we obtain the points of  $\tau_1^{\pm}(\omega, 1)$  and  $\tau_2^{\pm}(\omega, 1)$  by changing values of  $\omega$  with increment 0.0002 inside the interval  $[\bar{\omega}_1, \bar{\omega}_2]$  and then plot these points in the  $(\tau_1, \tau_2)$  plane. In **Figure 4A**, segment



 $L_2(1, 1)$  is illustrated as a real curve in the region above the diagonal and segment  $L_1(1, 1)$  is a dotted curve in the region below<sup>5</sup>. Two segments form an elliptical-shaped closed curve having the starting point  $S = (\tau_1^+(\bar{\omega}_1, 1), \tau_2^-(\bar{\omega}_1, 1))$  and the ending point  $E = (\tau_1^-(\bar{\omega}_2, 1), \tau_2^+(\bar{\omega}_2, 1))$  on the diagonal where, for i = 1, 2,

$$\tau_i^{\pm}(\bar{\omega}_1, 1) = \tau^s \simeq 8.478 \text{ and } \tau_i^{\pm}(\bar{\omega}_2, 1) = \tau^e \simeq 2.789.$$

Since the delays have two constraints,  $\tau_i \leq 5$  for i = 1, 2 and  $\tau_1 < \tau_2$  due to Assumptions 1 and 3, the feasible region should be above the diagonal line and subject to  $\tau_2 \leq 5$ . It is colored in yellow and further sub-divided into two sub-regions by the  $L_2(1, 1)$  segment that intersects two lines, one is the diagonal at  $(\tau^e, \tau^e)$  and the other is the horizontal line at  $\tau_2 = 5^6$ .

It is clear from the yellow region that stability is preserved for  $\tau_1$  and  $\tau_2$  such as  $0 < \tau_1 < \tau^e$  and  $\tau_1 < \tau_2 < \tau^e$ . When the varying pair  $(\tau_1, \tau_2)$  with  $\tau_2 > \tau^e$  crosses the  $L_2(1, 1)$ segment, we have the stability switching as the real part of at least one eigenvalue turns to be positive and then the stationary point loses stability. To see this, we perform a simulation by increasing the value of  $\tau_1$  along the dotted horizontal line at  $\tau_2 = 4$  that crosses the  $L_2(1, 1)$  curve at point  $(\tau_1^c, 4)$  with  $\tau_1^c \simeq 1.640^7$ . The numerical results are summarized in **Figure 4B** in which a bifurcation diagram with respect to  $\tau_1$  is depicted. For  $\tau_1 \leq \tau_1^c$ , a pair of  $(\tau_1, 4)$  stays within the stability region.

<sup>6</sup>The  $\tau_1$ -value of the intersection is obtained in the following way: first solving  $\tau_2^+(\omega, 1) = 5$  for  $\omega$  yields a solution  $\omega_m \simeq 0.560$  and then substituting it into  $\tau_1^-(\omega, 1)$  gives  $\tau_1^m \simeq 0.887$ . So the black segment depicted in the yellow region is defined for  $\omega \in [\omega_m, \omega_2]$ .

<sup>7</sup>It is possible to obtain this value by following the same procedure with which we obtained the value of  $\tau_1^m$  just above.



<sup>&</sup>lt;sup>5</sup>The segment  $L_i(k, n)$  shifts upward if *n* increases and rightward if *k* increases. For  $(k, n) \ge 2$  and (k, n) = 0, the switching segments are defined but located ouside of the region with  $\tau_i \le 5$  for i = 1, 2, so they are not depicted in **Figure 4A**.



So the stationary point  $p^* = f$  is stable and the corresponding part of the bifurcation diagram is a horizontal line at p(t) = f. Once  $\tau_1 > \tau_1^c$ , the pair is in the instability region above the  $L_2(1, 1)$  segment. The stationary point, therefore, loses stability and bifurcates to a limit cycle having two extreme values (i.e., maximum and minimum). The existence of the limit cycle is confirmed by the Hopf bifurcation theorem. **Figure 4B** further implies that the limit cycle becomes larger as  $\tau_1$  gets larger and the stability is never regained for  $\tau_1 \leq 4^8$ . The  $\tau_1$ -effect is summarized as follows:

Proposition 3. Stability of dynamic Equation (1) with respect to  $\tau_1$  depends on the selected value  $\bar{\tau}_2$  of  $\tau_2$ ;

- If  $\bar{\tau}_2 \leq \tau^e$ , then the steady state is stable for  $\tau_1 < \bar{\tau}_2$ ;
- If  $\overline{\tau}_2 > \tau^e$ , then there is the threshold value  $\tau_1^c$  such that the steady state is stable for  $\tau_1 < \tau_1^c$  and loses stability for  $\tau_1 \geq \tau_1^c$ .

where  $\tau_1^c$  is the  $\tau_1$ -point of the intersection of the switching curve and the horizontal line at  $\overline{\tau}_2$ .

Taking various values of  $\alpha_1$  from the interval  $[\alpha_1^0, 2/3]$  and constructing the stability switching curves, we detect the  $\tau_1$ -effect more. In **Figure 5A**, the switching curves with various values of  $\alpha_1$  are illustrated and their end-points on the diagonal are denoted by the green dots. The upper most blue curve has  $\alpha_1 = \alpha_1^0$  and the lower most blue curve has  $\alpha_1 = 2/3$  while the curve shifts downward as  $\alpha_1$  increases from  $\alpha_1^0$  to 2/3. As far as the curve is single-valued in  $\tau_2$ , we have essentially the same result on

<sup>8</sup>Mathematically, it might be possible to regain stability for  $\tau_1 > 4$ . However, economically it is not the case as  $\tau_1 \le \tau_2 \le 5$  is imposed by Assumptions 1 and 3.

the  $\tau_1$ -effect as in the case of  $\tau_1 = 6/15$ : fixing the value of  $\tau_2$  at some level and increasing the value of  $\tau_1$ , the stability is preserved for  $\tau_1$  below the threshold value  $\tau_1^c$  and it is lost for  $\tau_1 > \tau_1^c$ . As seen in **Figure 5A**, this threshold value becomes smaller as the value of  $\alpha_1$  becomes larger<sup>9</sup>, a larger value of  $\alpha_1$  reinforces the destabilizing  $\tau_1$ -effect.

# 3.2. Delay Effect II: $\tau_2$ -effect

Taking  $\tau_1$  fixed, we are next concerned with the  $\tau_2$ -effect, how changes in the value of  $\tau_2$  affect dynamics. We first return to **Figure 4A** in which both values of  $\alpha_1$  and  $\alpha_2$  are fixed at 6/15 and 1/3, respectively. The downward-sloping shape of the switching curve in the yellow region indicates that changing the value of  $\tau_2$  with constant  $\tau_1$  may generate similar effect to the  $\tau_1$ -effect. Let  $\tau_1^m$  denote the  $\tau_1$ -value of the intersection of the stability switching curve and the horizontal line at  $\tau_2 = 5^{10}$ . If  $\tau_1$  is selected to be less than  $\tau_1^m$ , then the steady state is stable for any  $\tau_2$  that should be above the diagonal and  $\leq 5$ . If  $\tau_1^m < \tau_1 < \tau^e$ , then the steady state is stable for  $\tau_2 < \tau_2^c$  and bifurcates to a limit cycle for  $\tau_2 > \tau_2^c$  where  $\tau_2^c$  is defined in the same way as  $\tau_1^c$  and depends on the selected value of  $\tau_1$ . One more case exists and is specific to the  $\tau_2$ -effect, namely, if  $\tau_1 > \tau^e$ , any feasible pair of  $\tau_1$  and  $\tau_2$  is in the unstable region so the steady state is unstable.

We now turn attention again to **Figure 5A** and consider the  $\tau_2$ -effect when  $\alpha_1$  is increased but  $\alpha_2$  is still fixed at 1/3. It is already seen that the stationary point is stable in the region left to the stability switching curve and increasing the value of  $\alpha_1$  shifts the switching curve leftward. The stability switching curve with  $\alpha = \alpha_0$  is not depicted as it is located outside **Figure 5A**. The shift of the curve implies that a larger value of  $\alpha_1$ 

<sup>&</sup>lt;sup>9</sup>This is not the case of the non-monotonics curve with  $\alpha_1 = 9/15$ .

<sup>&</sup>lt;sup>10</sup>" $\tau_1^m$ " is not labeled in **Figure 4A** to avoid figure congestion.

reinforces the destabilizing  $\tau_2$ -effect in the sense that it makes the stability region smaller. In addition, we have interesting dynamic phenomenon when the switching curve with  $\alpha_1$  being close to 2/3 has forward- and backward-bending segments. This shape suggests the stability loss and gain with respect to  $\tau_2$ . If the initial point of  $(\tau_1, \tau_2)$  is selected on the diagonal such that the stationary point is stable, then the vertical line standing at this value of  $\tau_1$  denoted as  $\tau_1^0 = 0.85$  crosses the stability switching curve three times at

$$\tau_2^a \simeq 1.014, \ \tau_2^b \simeq 1.687 \text{ and } \tau_2^c \simeq 3.359$$

as illustrated in **Figure 5A**. That is, the stability is lost at the first intersection and a limit cycle emerges. The cycle is getting larger and then smaller to merge the steady state at the value of the second intersection at which the stability is regained. Further increasing  $\tau_2$  crosses the stability switching curve again at which the stability is lost again. The corresponding bifurcation diagram is given in **Figure 5B** in which the stability losses and gains are shown. This  $\tau_2$ -effect is summarized as follows:

Proposition 4. With a value of  $\tau_1$  being close to 2/3, the increasing delay  $\tau_2$  has a double edge effect in a sense that it can destabilize the market price to generate a limit cycle, then stabilize it and later destabilize it again.

# 3.3. Delay Effect III: $\tau_3$ -effect

We now look into the  $\tau_3$ -effect, the effect caused by changing the value of  $\tau_3$ . To this end,  $\alpha_1 = (6 - 0.1)/12$  is taken and dynamics with three different values of  $\tau_3$ ,  $\tau_3 = 5$ ,  $\tau_3 = 6$ , and  $\tau_3 = 7$ , are considered. The corresponding switching curves are illustrated in different colors in **Figure 6A** (that is, a blue curve with  $\tau_3 = 5$ , red curves with  $\tau_3 = 6$ , and black curves with  $\tau_3 =$ 7). As  $\tau_3$  increases, it is observed first that the feasible regions of  $\tau_1$  and  $\tau_2$  increase since the upper bound of  $\tau_2$  becomes larger and second that the downward-sloping switching curve shifts upward, which implies an enlargement of the stability region. It is further observed that island-shaped switching curves emerge for  $\tau_3 = 6$  and  $\tau_3 = 7$  and the steady state is unstable inside them. To detect the reason why the island-shaped region is born under larger values of  $\tau_3$ , we draw the green  $f_1(\omega) = f_2(\omega)$  curve and the red  $f_d(\omega)$  curve for  $\tau_3 = 7$  in **Figure 6B**. It is seen that  $f_1(\omega) = f_2(\omega) > 0$  for  $\omega > 0$  and the larger  $\tau_3$  value makes the second hump of the  $f_d(\omega)$  curve high enough to cross over the horizontal line. As a result, the red curve intersects the horizontal line four times at

$$\omega_1 \simeq 0.104, \omega_2 \simeq 0.562, \omega_3 \simeq 1.059 \text{ and } \omega_4 \simeq 1.256.$$

In consequence, the triangle conditions (5) hold in the two intervals,  $[\omega_1, \omega_2]$  and  $[\omega_3, \omega_4]$ . It is verified that the switching curve defined on  $[\omega_1, \omega_2]$  is downward-sloping and the one on  $[\omega_3, \omega_4]$  forms an island shape. Both curves are colored in black. The existence of the smaller red islands indicates that the  $f_d(\omega)$ curve also intersects the horizontal axis four times even when  $\tau_3 = 6$ . A critical value of  $\tau_3$  for the birth of the island is somewhere between  $\tau_3 = 5$  and  $\tau_3 = 6$ , for which the maximum of the second hump of  $f_d(\omega)$  becomes zero.

Proposition 5. As far as  $\tau_3$  is relatively small, increasing the value of  $\tau_3$  stabilizes the steady state by shifting the downward-sloping switching curve upward whereas for a relatively larger value, increasing value of  $\tau_3$  additionally generates an island-shaped switching curve within which the steady state is unstable.

To see the  $\tau_3$ -effect on dynamics, we perform two simulations. In the first simulation, we increase the value of  $\tau_1$  along the horizontal line at  $\tau_2 = \tau_2^0 (\simeq 4.356)$ . This line successively





crosses the stability switching curves, the blue curve at  $\tau_1^a (\simeq 0.972)$ , the red curve at  $\tau_1^b (\simeq 1.372)$  and the black curve at  $\tau_1^c (= 2)$ . Although, these bifurcating values are denoted as black dots, they are not labeled in **Figure 6A** to avoid confusion. The corresponding bifurcation diagrams are illustrated in **Figure 7A** in which the black curve is depicted first, the red curve is then put on it and finally the blue curve is further placed upon. Each diagram has a qualitatively similar shape as the one shown in **Figure 4B**, that is, the steady state loses stability at the bifurcation value and a limit cycle emerges for a larger value of  $\tau_1$ . Thus larger  $\tau_3$  value does not alter qualitative aspects because the bifurcation value of  $\tau_1$  becomes larger as  $\tau_3$  increases. Increasing value of  $\tau_3$  enhances the stabilizing  $\tau_1$ -effect by delaying the loss of stability.

In the second simulation, we shift the emphasis away from the  $\tau_1$ -effect to the  $\tau_2$ -effect. Fixing  $\tau_3 = 7$  and starting at point (2, 2) on the diagonal, we increase the value of  $\tau_2$  from 2 to 7 along the vertical real line at  $\tau_1 = \tau_1^0 (= 2)$  in **Figure 6A**. The line crosses the lower black island-shaped curve at  $\tau_2^a (\simeq 2.563)$ , the black downward sloping curve at  $\tau_2^b (\simeq 4.356)$  and then upper black circle-shaped curve at  $\tau_2^c (\simeq 6.283)$ . These threshold values are not labelled in **Figure 6A** by the same reason as before. The simulation results are given in **Figure 7B** in which stability gain and loss occur in the following way:

(i) we initially have a unstable steady state (in other words, a stable limit cycle) when the initial point is selected on the diagonal in Figure 6A<sup>11</sup>, that is inside the half black circle. The real part of one eigenvalue is positive.

- (ii) the limit cycle becomes smaller as  $\tau_2$  increases from 2 to  $\tau_2^a$  and merges with the steady state for  $\tau_2 = \tau_2^a$  at which instability is switched to stability. The positive real part turns to be negative when the delay  $\tau_2$  crosses the imaginary axis from left to right.
- (iii) for  $\tau_2^a < \tau_2 < \tau_2^b$ , the model has a stable steady state and stability loss occurs for  $\tau_2 = \tau_2^b$ . The real parts of at least one eigenvalue turn to be positive again when the delay  $\tau_2$  crosses the imaginary axis.
- (iv) for  $\tau_2 > \tau_2^b$ , **Figure 7B** suggests the existence of stable limit cycles and the instability of the steady state. Although further increasing  $\tau_2$  intersects the island switching curve at point  $(2, \tau_2^c)$  at which the real part of another eigenvalue changes sign. However no stability switch occurs, since there is already at least one other eigenvalue with positive real part.

Proposition 6. Increasing the value of  $\tau_3$  does not affect essentially dynamics with respect to  $\tau_1$  and can affect the dynamics with respect to  $\tau_2$  as it can generate multiple stability switching.

## 4. CONCLUDING REMARKS

This paper constructs a heterogeneous agent model with three fixed delays and considers its dynamic behavior both analytically and numerically. The dependence of the delay effects on the changes in the lengths of the delays is also studied. The common result is that a longer delay can destabilize the market and can give rise to cyclic oscillations around the equilibrium (i.e., fundamental) price. This result clarifies the instability condition and thus complements the numerical

<sup>&</sup>lt;sup>11</sup>More precisely, we select a constant initial function, g(t) = 2 for  $t \le 0$ .

study of Dibeh [9]. In addition, it is found that under multiple delays, stability loss and gain repeatedly occur as the length of a delay increases. He and Zheng [6] observed this phenomenon in a financial market model with continuously distributed time delay with finite memory length that involved infinitely many past price observations and called it the double edge effect. Our model is based on one, two or three past observations. It is thus shown in this paper that the same effect can occur even under finite delay information. In our future research the model of this paper will be extended to involve two risky assets and the stability and instability of the equilibrium will be examined in the extended models.

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# **AUTHOR CONTRIBUTIONS**

AM: Performs simulations; FS: Constructs the model.

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